

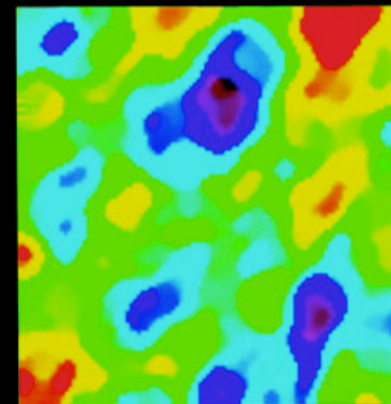
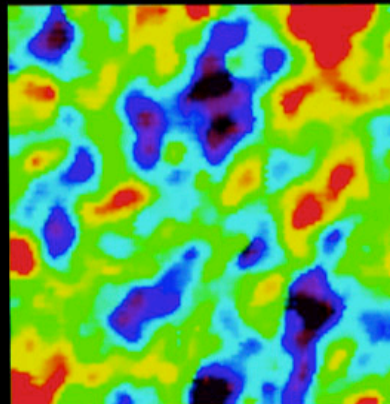
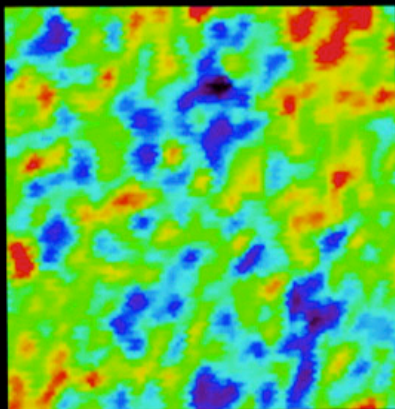
Precision Measurement of the Mean Curvature

Lloyd Knox

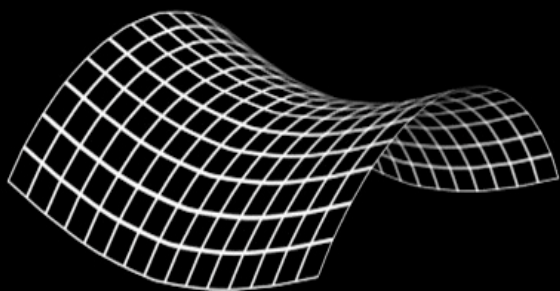
University of California, Davis

astro-ph/0503405

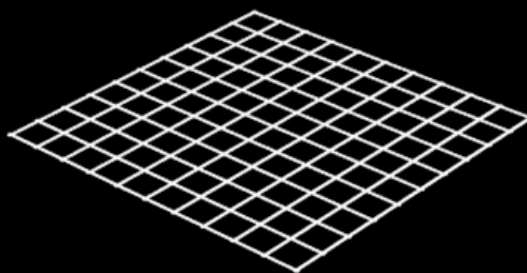
GEOMETRY OF THE UNIVERSE



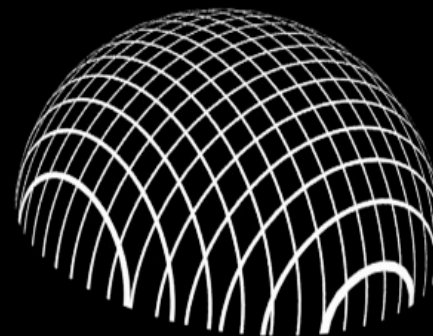
Physical size of typical hot/cold spot can be calculated. How this projects into angular size depends on curvature.



OPEN



FLAT

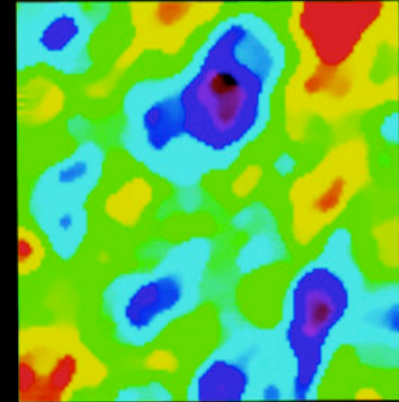
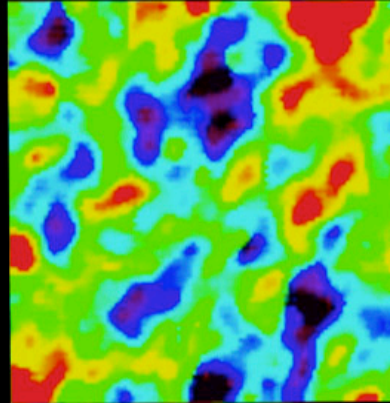
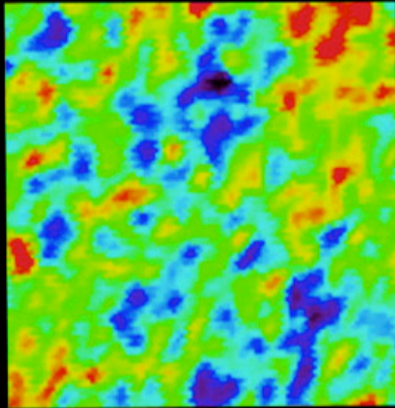


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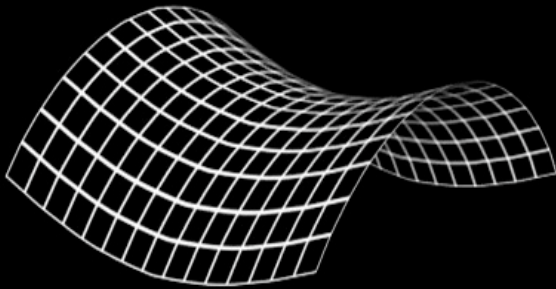
“Weighing the Universe with the CMB”

Jungman et al., Phys.Rev.Lett. 76, 1007 (1996).

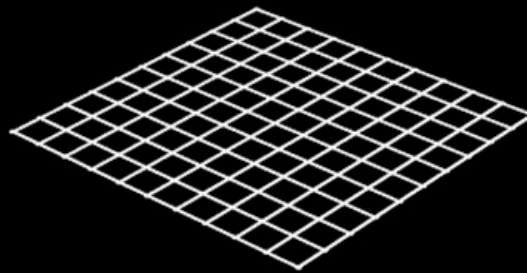
GEOMETRY OF THE UNIVERSE



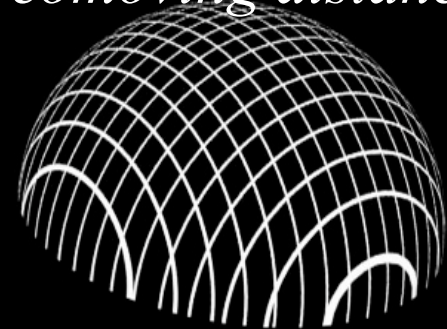
Physical size of typical hot/cold spot can be calculated. How this projects into angular size depends on curvature and *comoving distance*.



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So how has $\Omega_{\text{tot}} \approx 1$ been inferred from CMB data? Answer:
 $\partial D_A / \partial \Omega_{\text{tot}} \approx 5 \partial D_A / \partial \Omega_\Lambda \rightarrow$ If $\sigma(\Omega_\Lambda) = 0.5$ then $\sigma(\Omega_{\text{tot}}) = 0.1$

Outline

- Importance of Mean Curvature Measurement
- Dark Energy / Curvature Degeneracy
- A Straightforward Solution
- Standard Candles and Standard Rulers
- Conclusion

Why Measure Mean Curvature?

- Robust Prediction of Inflation

$$\langle \rho \rangle / \rho_c = 1 \pm 10^{-60}$$

$$\langle \rho \rangle_H / \rho_c = 1 \pm 10^{-5} \quad (\text{averaged over Hubble patch})$$

- Probe of Fluctuations on Super-horizon Scales

How Well Is it Known Already?

- If we assume the dark energy is a cosmological constant, the SDSS baryon oscillation detection combined with CMB data gives a very impressive constraint of $\Omega_{\text{tot}} = 1.01 \pm 0.009$

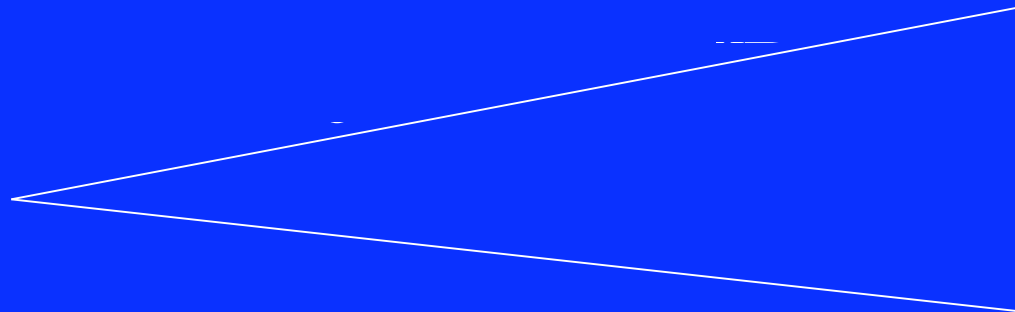
But we don't know that the dark energy is a cosmological constant. You may have noticed there's a minor effort underway to investigate the nature of the dark energy. If we allow w not equal to -1 then this constraint weakens considerably.

Dark Energy / Curvature Degeneracy

The comoving size of the sound horizon depends on matter density and baryon density, which can be inferred from CMB acoustic peak morphology, and thereby calibrated.

Measure θ and infer D_A

$$\Delta D_A / D_A = \Delta r_s / r_s - 0.25 \Delta \rho_m / \rho_m$$



Standard ruler



But D_A depends on both curvature and matter content \rightarrow degeneracy

$\Omega_\Lambda - \Omega_k$ degeneracy: Eisenstein et al. (1998), Efstathiou and Bond (1999)

Dark Energy / Curvature Degeneracy

$$ds^2 = dt^2 - a^2(t) [dr^2/(1-kr^2) + r^2(d\theta^2 + \sin^2\theta d\phi^2)]$$

From line element, $D_A = r$, and comoving distance from origin to r is

$$l = s_0 \int_0^r dr' / (1-kr'^2)^{1/2}$$

Solving for r to lowest order in k we have

$$D_A = r = l + kl^3/6$$

If we knew l and D_A we could solve for k . But we don't know l .

Instead, we can calculate the comoving distance traveled by a photon that suffers a redshift, z :

$$l(z) = s_0 \int_0^z dz' / H(z') \quad \text{where} \quad H^2(z) = 8\pi G\rho(z)/3 - k/a^2$$

Precision Determination of Mean Curvature



Measure D_{OL} (with CMB) and D_{OM} (e.g., baryon oscillations)

Calculate l_{ML} (given ρ_m from CMB)*

In absence of curvature, $D_{OL} - (D_{OM} + l_{ML}) = 0$

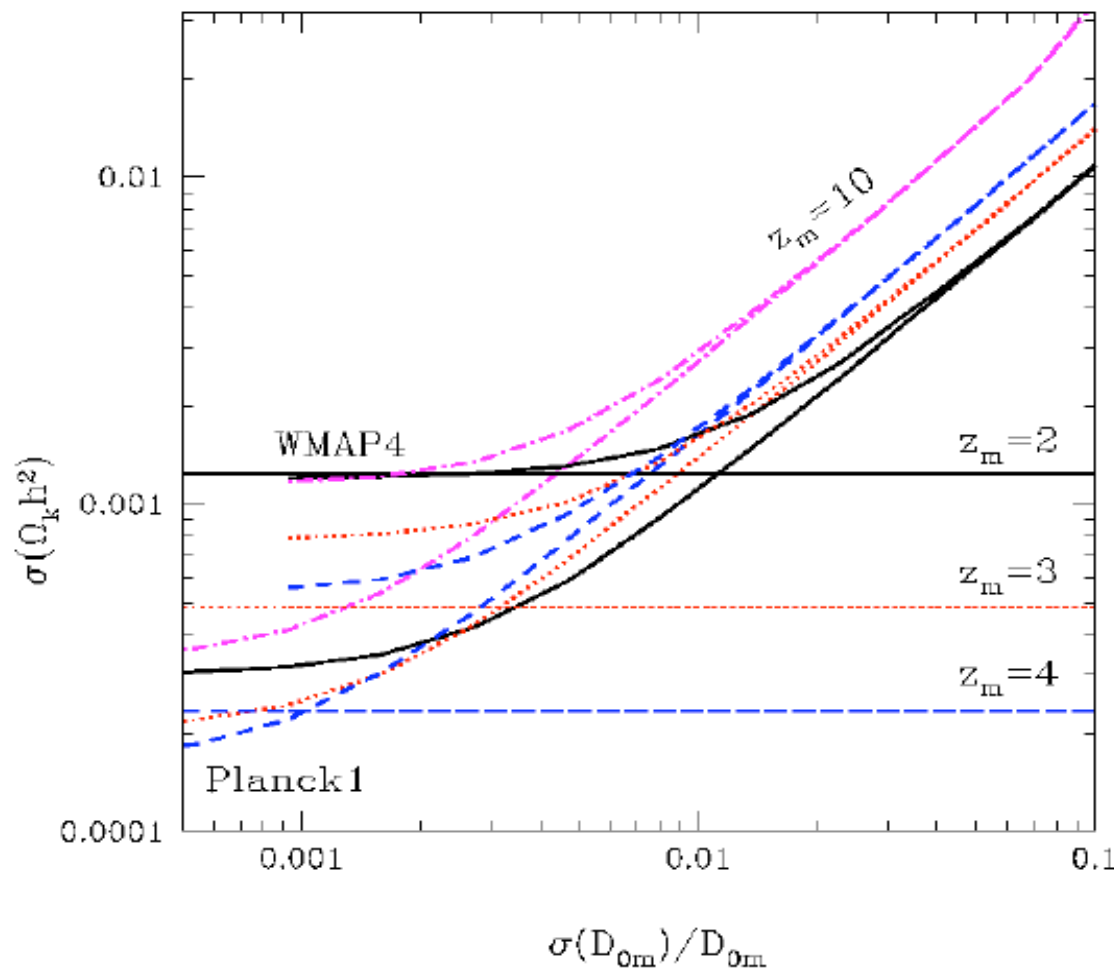
More generally (for $|\Omega_k| \ll 1$):

$$D_{OL} - (D_{OM} + l_{ML}) = \Omega_k H_0^2 (D_{OL}^3 - D_{OM}^3) / 6$$

*Note: l_{ML} is the comoving distance, equal to angular diameter distance D_{ML} if $\Omega_k = 0$.

Error on Curvature Given Error on D_{OM}

$$\Omega_k h^2 = (h/H_0)^2 (D_{OL} - (D_{OM} + l_{ML})) / (D_{OL}^3 - D_{OM}^3)$$



Horizontal lines: bias in method due to dark energy at $z > z_m$

Other lines: error in $\Omega_k h^2$ due to CMB errors on $\Omega_b h^2$ and $\Omega_m h^2$ as well as D_{OM} measurement error.

In limit of perfect D_{OM} , errors in D_{OL} and l_{ML} (due to error in ρ_m) partially cancel.

10^{-5} is difficult!

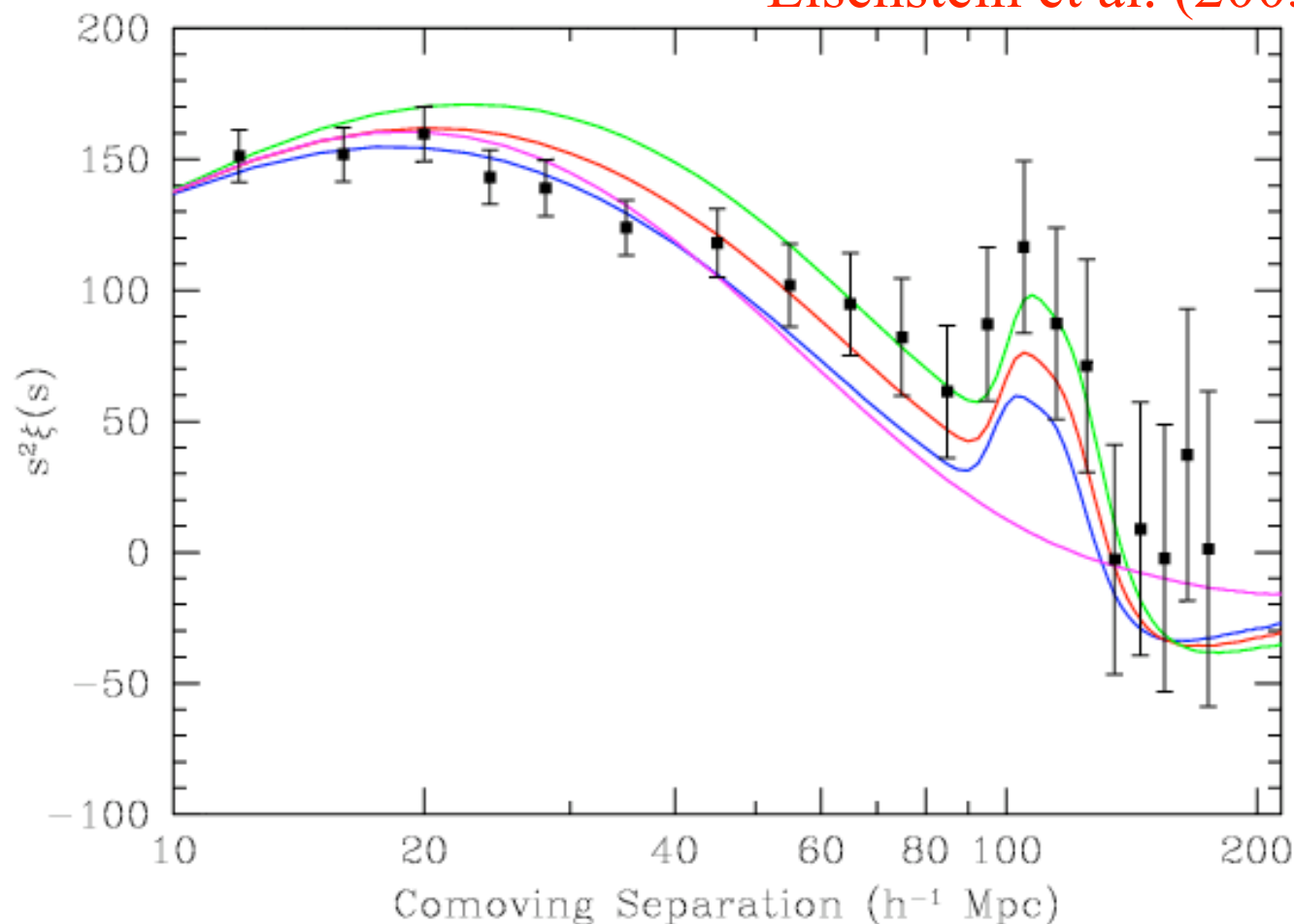
Measuring D_{OM}

- Standard Candles
 - SNeIa, GRB?, ??
- Standard Rulers in Matter Power Spectrum
 - sound horizon at last-scattering:
 $r_s \rightarrow \text{measure } D_A/r_s$
 - particle horizon at matter-radiation equality: $\quad \quad \quad /$
 $1/\rho_{m,0} \rightarrow \text{measure } D_A \rho_{m,0} \text{ (or } D_A \omega_m \text{)}.$

$$(\omega_m = \rho_{m,0}/\rho_c h^2)$$

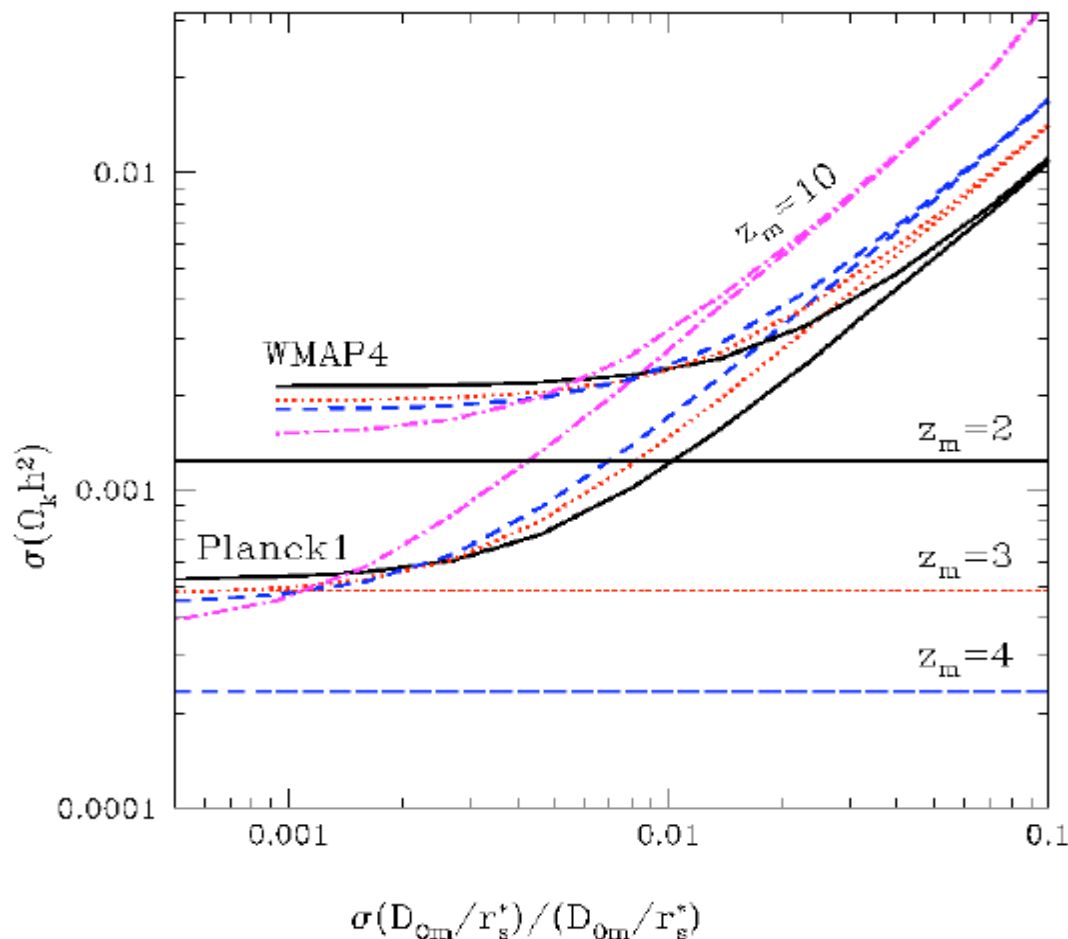
Correlation Function of SDSS Luminous Red Galaxies

Eisenstein et al. (2005)



Curvature Error Given Error on D_{OM}/r_s

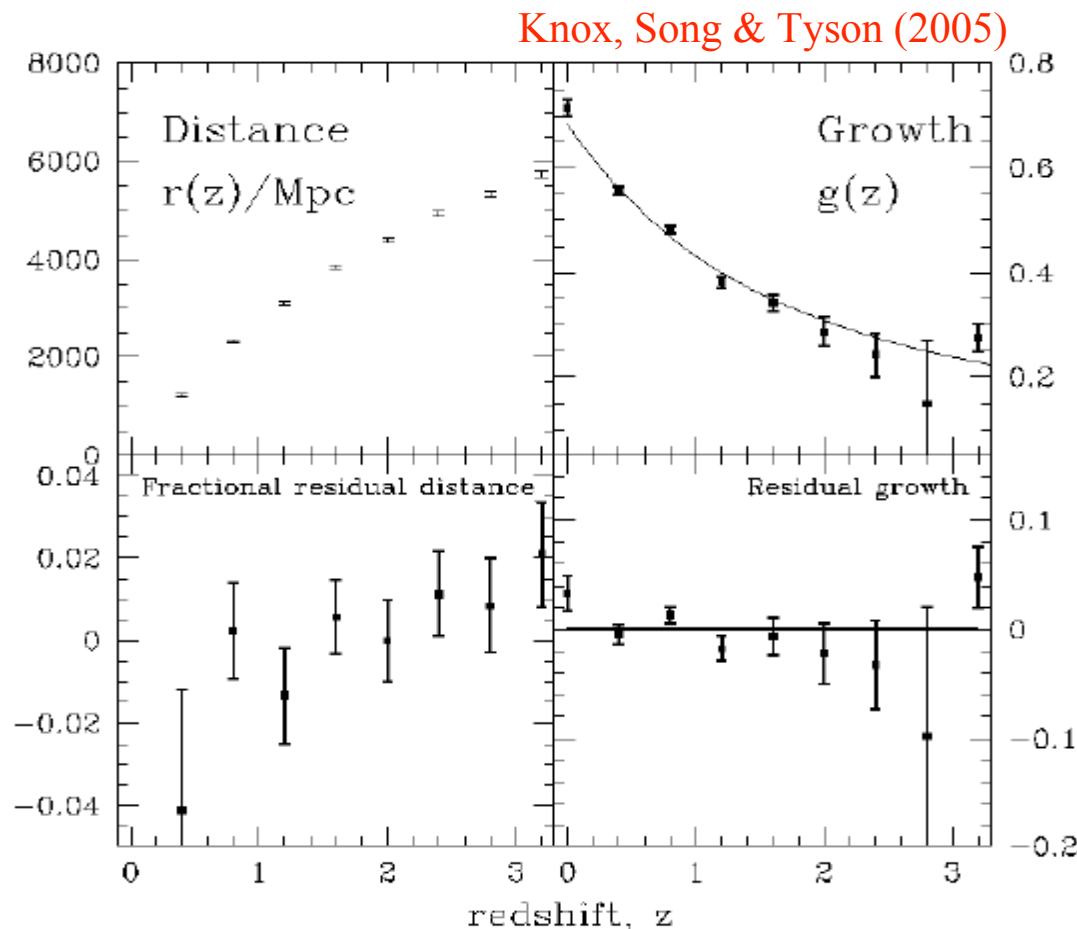
$$\Omega_k h^2 = 6(h/H_0)^2 r_s^{-2} (D_{\text{OL}}/r_s - (D_{\text{OM}}/r_s + l_{\text{ML}}/r_s)) / ((D_{\text{OL}}/r_s)^3 - (D_{\text{OM}}/r_s)^3)$$



No significant error in $D_{\text{OL}}/r_s (=1/\theta_s)$.

In limit of perfect D_{OM} , error is entirely from l_{ML} error.

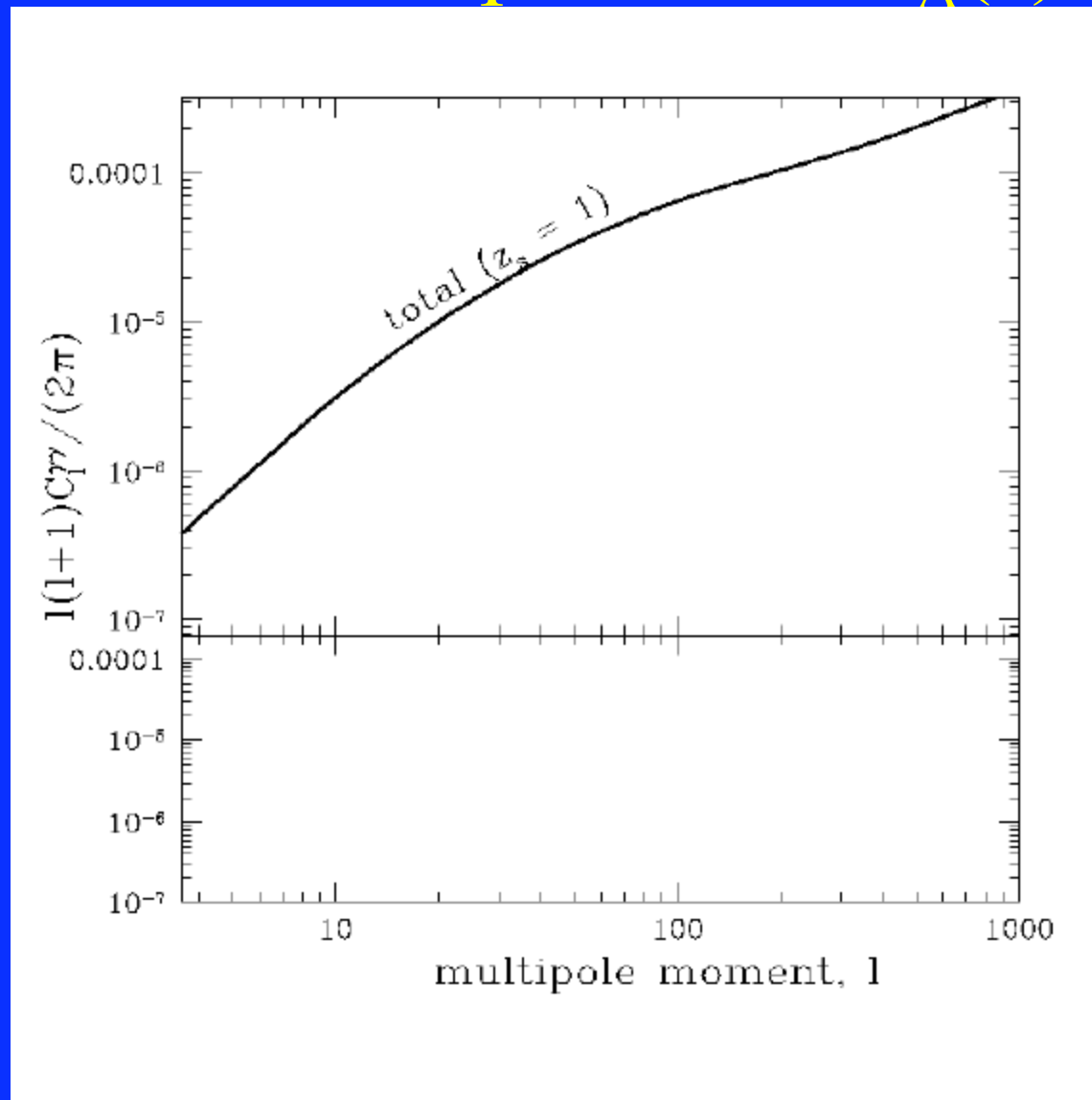
Distance (and growth) reconstructed from LSST WL survey + Planck



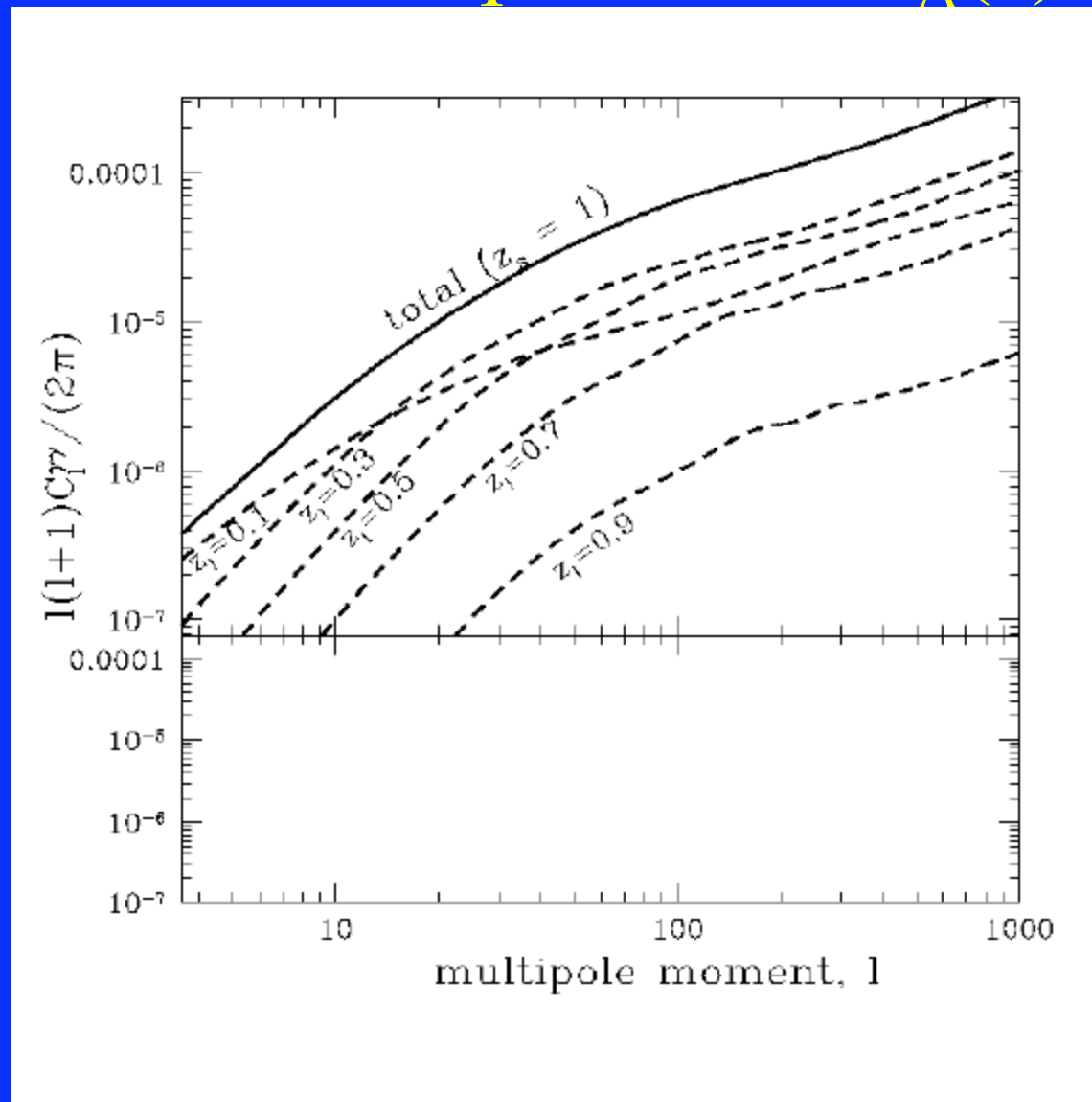
With the parameters of the high- z Universe pinned down by Planck, only thing left to measure is $g(z)$ and $D_A(z)$ (here called $r(z)$) in the dark energy-dominated era. They can both be reconstructed from tomographic cosmic shear data.

D.E. constraints come almost entirely from $D_A(z)$ constraints (Simpson & Bridle '04, KST05) .

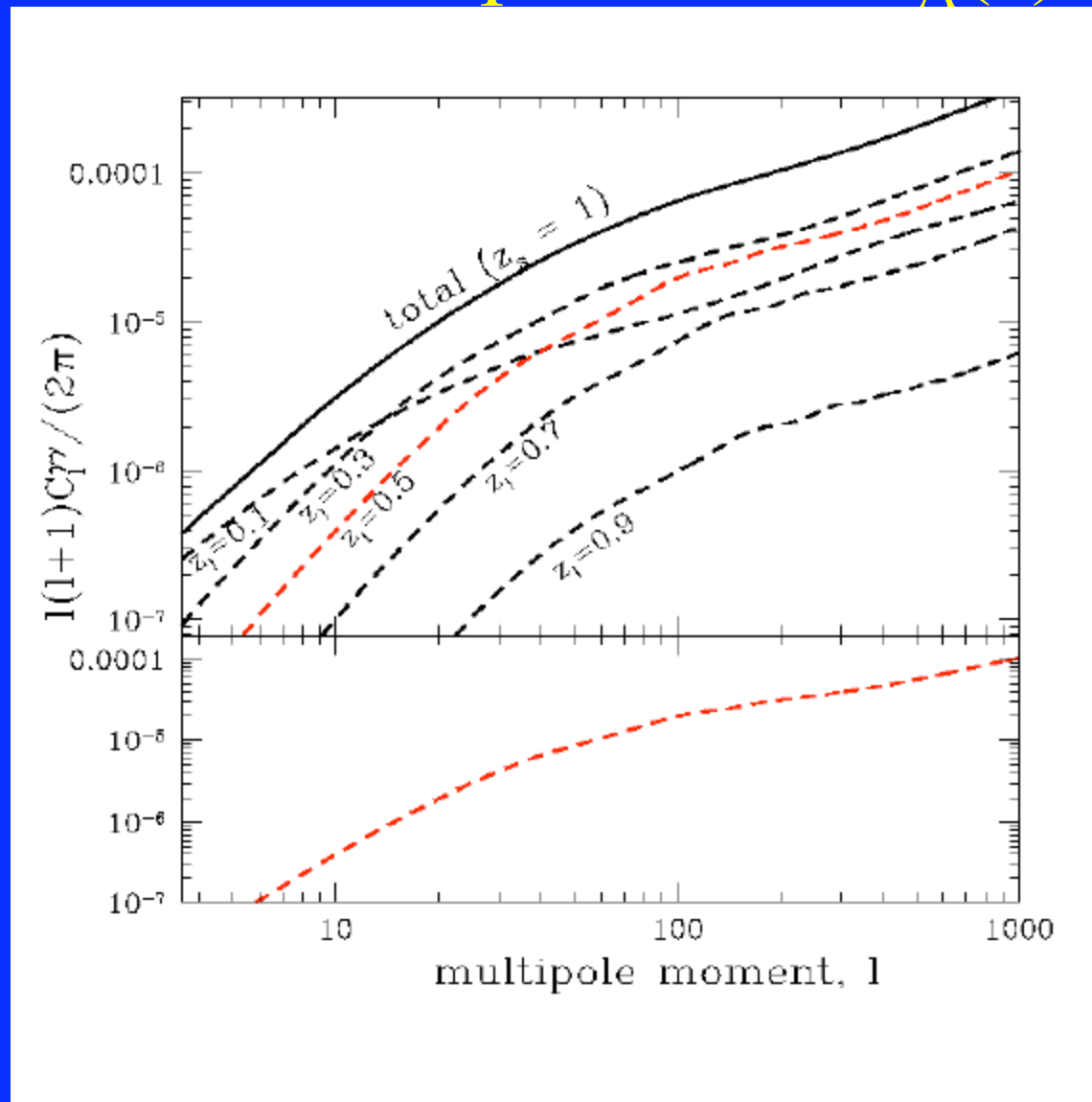
Dependence of Shear power on $D_A(z)$ and $g(z)$



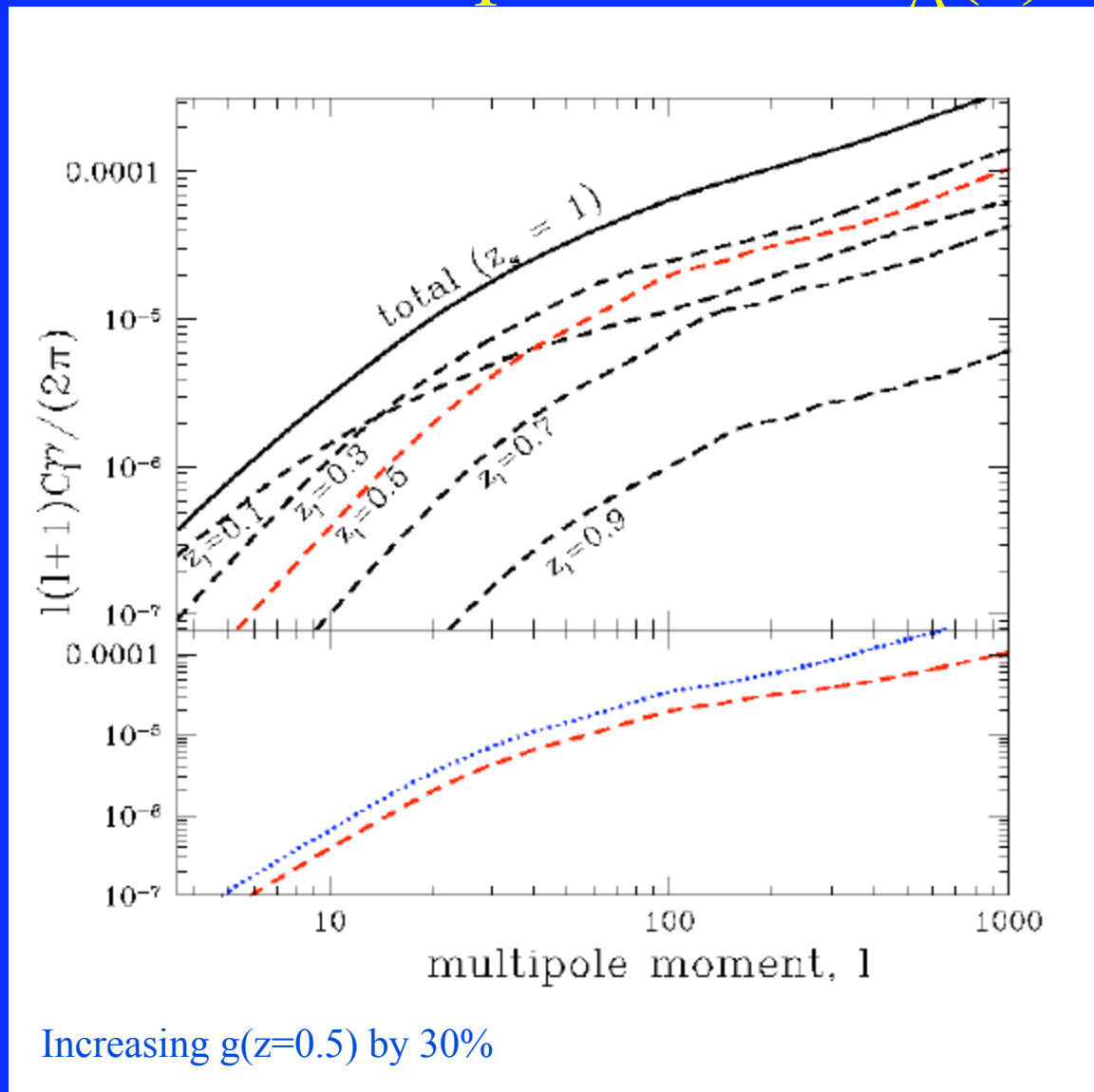
Dependence of Shear power on $D_A(z)$ and $g(z)$



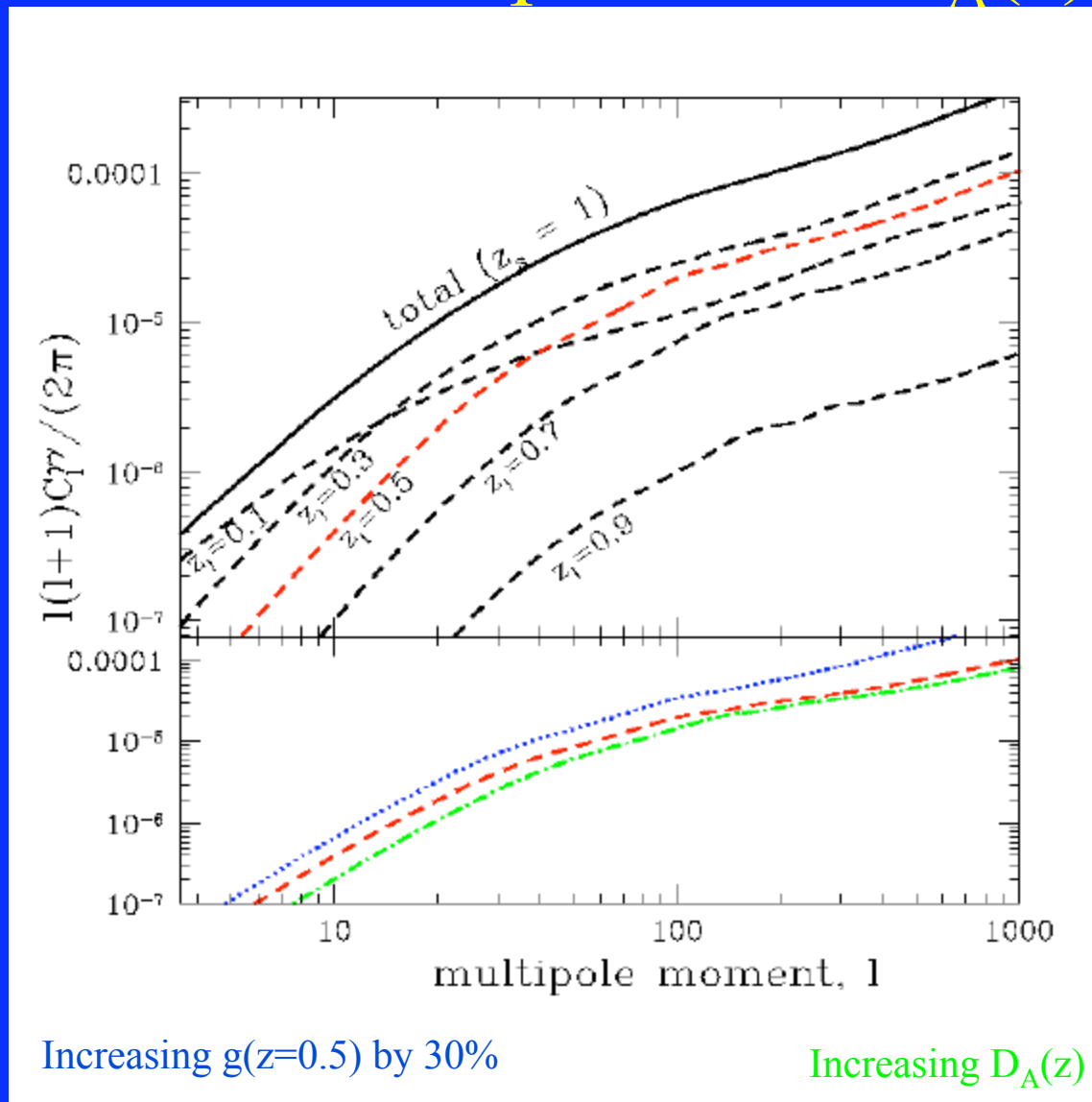
Dependence of Shear power on $D_A(z)$ and $g(z)$



Dependence of Shear power on $D_A(z)$ and $g(z)$

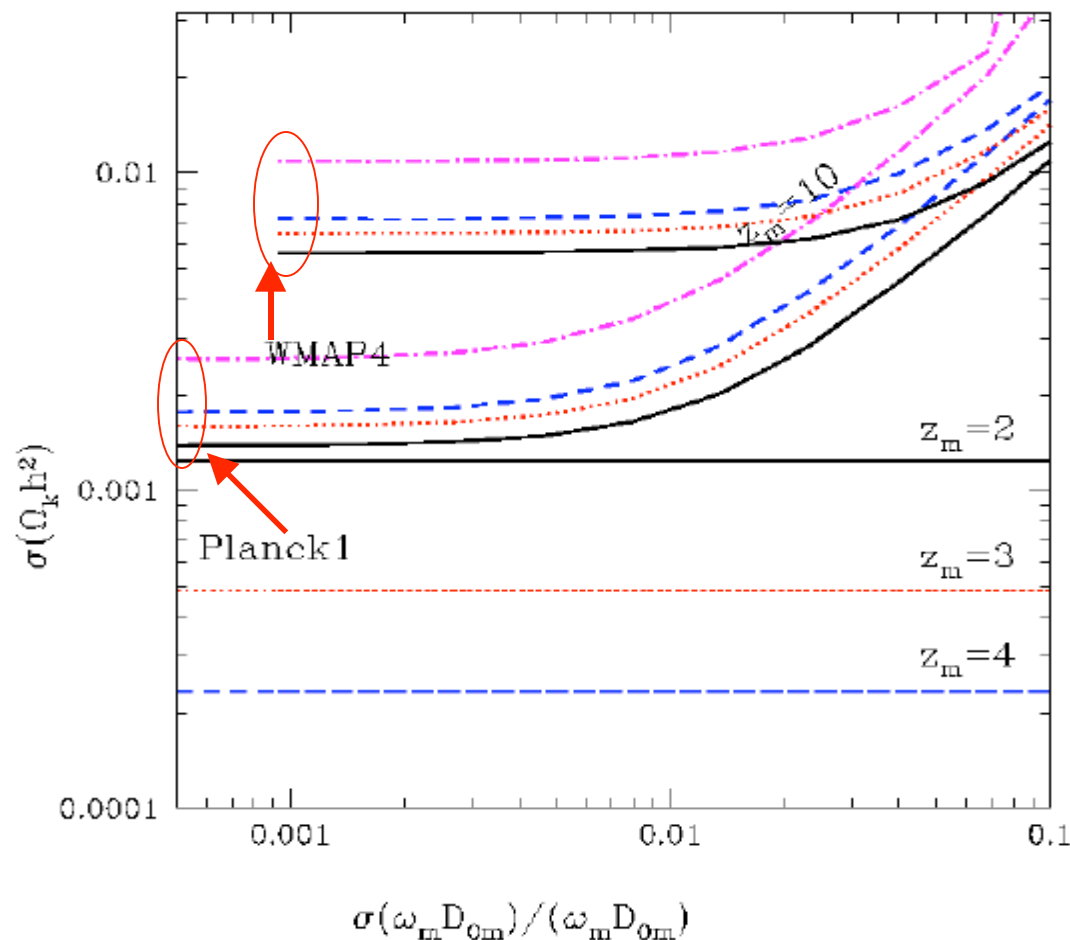


Dependence of Shear power on $D_A(z)$ and $g(z)$



Curvature Error Given Error on $D_{OM}\omega_m$

$$\Omega_k h^2 = 6(h/H_0)^2 \omega_m^2 (D_{OL}\omega_m - (D_{OM}\omega_m + l_{ML}\omega_m)) / ((D_{OL}\omega_m)^3 - (D_{OM}\omega_m)^3)$$



Limit of perfect $D_{OM} \omega_m$:
Cancellation no longer as
good between $D_{OL} \omega_m$ and
 $l_{ML} \omega_m$

We do significantly worse
here than in pure distance
measurement case or in
baryon oscillation case.

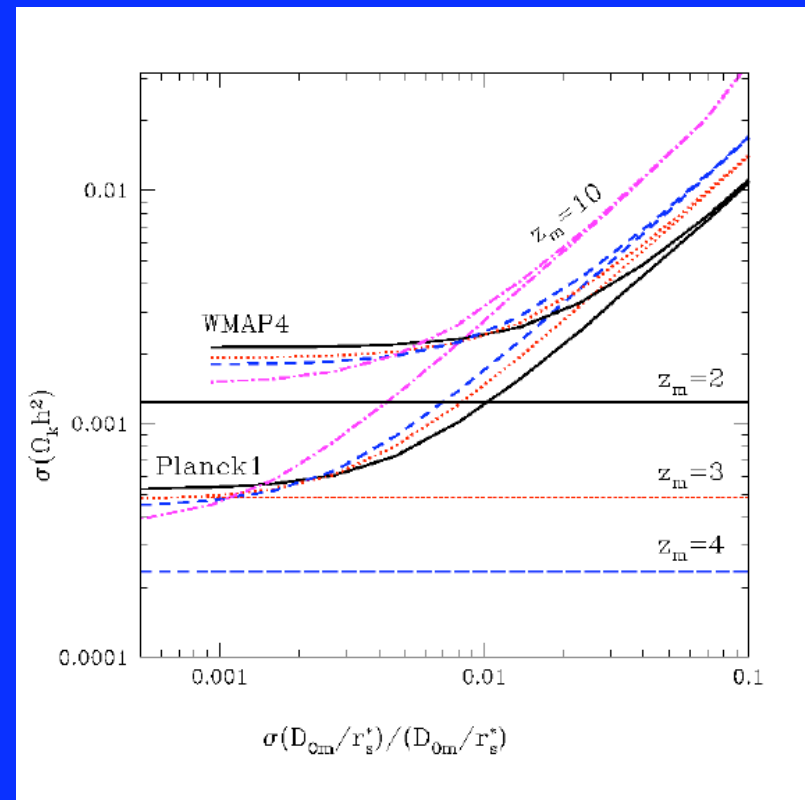
Note on Robustness of $\Omega_k h^2$ from BAO

- CMB acoustic peak morphology affected by evolution of gravitational potentials \rightarrow constrains ρ_m/ρ_{rad} and therefore ρ_m *if we know radiation content*.
- Independent of radiation content CMB robustly constrains $\rho_m^{1/2} r_s$.
- Since BAO constrain D_A/r_s and we know $\rho_m^{1/2} r_s$ we actually learn $D_A \rho_m^{1/2}$ (Eisenstein & White (2004))
- $\Omega_k h^2 / D_{\text{OL}} \omega_m^{1/2} - \underbrace{(D_{\text{OM}} \omega_m^{1/2} + 1_{\text{ML}} \omega_m^{1/2})}$

Has no dependence on
cosmological parameters!

What would a detection at 10^{-3} level possibly mean?

- Inflation did not happen (but then what did that leaves small curvature?)
- Inflation occurred and ended with bubble nucleation followed by ~ 60 e-folds of slow-roll. [Very fine-tuned!]
- Extra fluctuation power on super-horizon scales.



Another Way to Measure Mean Curvature

Bernstein (2005)

It's always true that

$$r_{AC} - (r_{AB} + r_{BC}) = 0$$

where A is the origin.

It's also true that $D_{AC} = r_{AC}$ and $D_{AB} = r_{AB}$,

but D_{BC} is *not* equal to r_{BC}

In fact,

$$D_{AC} - (D_{AB} + D_{BC}) / \Omega_k$$

WL is sensitive to all three distances. BAO can help.

Summary

- Zero mean curvature is a robust prediction of inflation worth rigorous checking.
- Uncertainty about dark energy limits our current knowledge of the mean curvature.
- Measurement of distances into the matter-dominated era will greatly reduce the dark energy model-dependence of any curvature determination.

• $\Omega_{\text{total}}=1$ ✓

WMAP

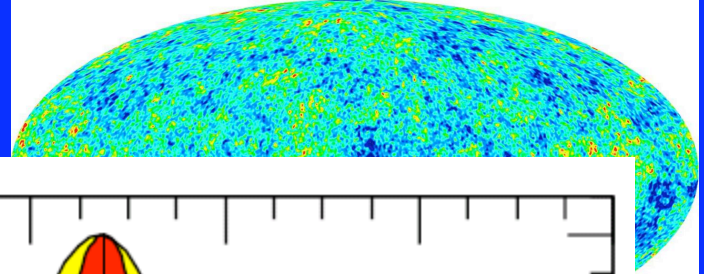
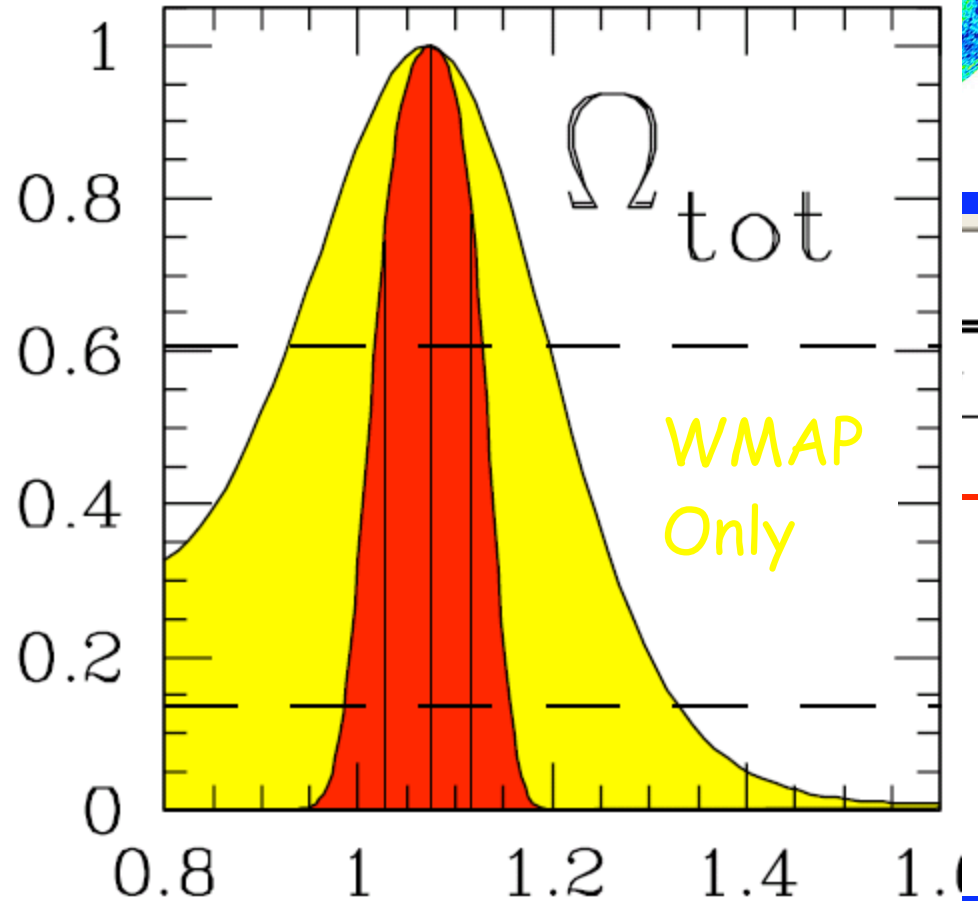


Table 3. "B

Description
Total density
Equation of state of quintessence
Dark energy density
Baryon density
Baryon density
Baryon density (cm^{-3})
Matter density
Matter density
Light neutrino density



Tegmark et al

astro-ph/0310723

WMAP
+ SDSS

Bennett et al Feb 11 '03